Challenges of 4th Generation Light Sources

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Abstract
We show that using existing electron gun technology and a high energy linac like the one at SLAC, it is possible to build a Free Electron Laser operating around the 4 nm water window. A modest improvement in the gun performance would further allow to extend the FEL to the 0.1 nm region. ..with a brightness many order of magnitude above that of any synchrotron radiation source, …in the multi gigawatt region and subpicosecond pulse length. ..
X-ray FELs, generate coherent electromagnetic radiation at \(~1\text{Å}, 1 \text{ fs}\), atomic phenomena characteristic length and time scales, with high peak power and brightness, giving us new, unprecedented capabilities to study the structure and dynamics of atomic and molecular systems, for biology, physics and chemistry, complex materials, matter under extreme conditions.

The X-ray FEL radiation: \(~10^9\) photons in a coherence volume compared to less than 1 for spontaneous radiation. This property allows studies of coherent diffraction imaging, multi-photons excitations, non-linear X-ray experiments, largely unexplored areas of science.
Outline

- Present status
- Next steps and next generations
- The Terawatt X-ray FEL
- Electron beam physics challenges
- Laser/plasma/wake-fields accelerators
- Conclusions

Disclaimer: X-ray FELs are today a very active field of research, development and construction. I cannot even try to cover it all in 30 minutes. I have made some choices reflecting my own bias. FLASH, SACLA, Fermi@Elettra have already been discussed in papers presented at this conference and I will mostly refer to LCLS data.
Present status: LCLS, long pulse case*
*P. Emma et al., Nature Photonics, DOI: 10.11038 2010.176

<table>
<thead>
<tr>
<th></th>
<th>Electrons</th>
<th>Xrays</th>
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<tbody>
<tr>
<td>Charge/bunch, nC</td>
<td>0.25</td>
<td>0.25</td>
</tr>
<tr>
<td>Beam energy, GeV</td>
<td>13.6</td>
<td>3.5–6.7</td>
</tr>
<tr>
<td>Slice emittance (injected), µm</td>
<td>0.4</td>
<td>0.4</td>
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<tr>
<td>Projected emittance, µm</td>
<td>0.5–1.2</td>
<td>0.5–1.6</td>
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<tr>
<td>Peak current, kA</td>
<td>2.5–3.5</td>
<td>0.5–3.5</td>
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<tr>
<td>Radiation wavelength, Å</td>
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<td>6–22</td>
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<tr>
<td>FEL gain length, m</td>
<td>3.5</td>
<td>1.5</td>
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<tr>
<td>Photons per pulse x 10^{12}</td>
<td>1.0–2.3</td>
<td>10–20</td>
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<tr>
<td>Peak X-ray power, GW</td>
<td>15–40</td>
<td>3–35</td>
</tr>
<tr>
<td>Pulse length (FWHM), fs</td>
<td>70–100</td>
<td>70–500</td>
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<tr>
<td>Bandwidth (FWHM), %</td>
<td>0.2–0.5</td>
<td>0.2–1.0</td>
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</table>
Low charge LCLS, 20pC


Peak power along the undulator and snapshot of typical 2fs FEL pulse at 100 m. Notice the characteristic SASE spikes.

800 eV. SASE Intensity fluctuations, corresponding to about 5 longitudinal modes. J. Wu et al., Proc. 2010 FEL Conf.
A high gain FEL is characterized by the FEL parameter, $\rho$, [R. Bonifacio, C. Pellegrini, and L.M. Narducci, Opt. Commun., 50, 373 (1984)] giving:

1. the exponential growth rate, $P = P_0 \exp(z/L_G)$, where $L_G \sim \lambda_U/4\pi\rho$

2. The saturation power $P_{sat} = \rho I_{beam} E_{beam}$

$$\rho = \left( \frac{K \Omega_p}{4 \omega_U} \right)^{2/3}, \quad \Omega_p = \left( \frac{4\pi r_e c^2}{\gamma^3 n_e} \right)^{1/2}, \quad \omega_U = \frac{2\pi c}{\lambda_U}$$

For LCLS: $\lambda = 1.5\AA$, $I_{beam} \sim 3kA$, $E_{beam} \sim 14\text{ GeV}$, $\rho \sim 5\times10^{-4}$, $P_{sat} \sim 20\text{ GW}$, $N_{coh} \sim 10^{12}$ in 100fs. Typically $\rho \sim 10^{-3}-10^{-4}$ for soft and hard X-rays

The number of coherent photons/pulse scales almost linearly with the pulse duration $\sim 10^{12}$ at 100 fs, $10^{11}$ at 10fs for 8keV- and inversely with the photon energy.
Remarks on 4th generation

• FLASH and LCLS have demonstrated outstanding capabilities, increasing by 7 to 10 orders of magnitude the photon peak brightness.

• The LCLS X-ray pulse duration and intensity can be changed from about 100 to a few femtosecond and $10^{13}$ to $10^{11}$ photons/pulse, over the wavelength range of 2.2 to 0.12 nm, by varying the electron bunch charge from 250 to 20 pC. **The X-ray pulse can be optimized for the experiment, not possible in storage ring sources.**

• Theory, simulations and experimental results are mostly in good agreement. Simulations tools have been developed and benchmarked to evaluate the electron beam properties and the X-ray pulse characteristics, from the electron source to the undulator exit. These tools can be used to design new, advanced FELs.

• We know that we can generate high energy electron beams with phase space density larger than what we expected until recently.
Electron beam (The Lasing Medium) physics challenges

a) Beam noise and density modulation from electron gun generates large coherent transition radiation and coherent synchrotron radiation, limiting the bunch compression, beam peak current and the FEL gain and power. An effect not predicted in the design phase and not yet fully understood.

b) Limitations to electron bunch compression and peak current, non flat bunch current profile after compression (the double horn current profile), have an effect on the FEL performance and design.

Understanding and controlling these effects, generating a better lasing medium, leads to more powerful and more compact X-ray FEL.
Magnetic bunch compression

LCLS and other 4th generation FELs use magnetic compression with two chicanes as shown for LCLS.

Coherent synchrotron radiation (CSR) and space charge effects during the compression can increase the beam emittance and distort the phase space.

Magnetic bunch compression

LCLS measurements of emittance and compression at 250 pC. Coherent synchrotron radiation (CSR) blows up the emittance for large compression. The effect is larger at large charge.
Laser heater

A “Laser Heater” has been developed and used in LCLS to increase the beam energy spread, and control these effects, yielding larger X-ray pulse power at 250 pC.


Laser heater decreases gain length.
Longitudinal phase space measurements after the second bunch compressor, and before the final acceleration. Under-compression and over-compression phase space and current profile. Beam energy $\sim 4$ GeV. The bunch head to the left. Y. Ding et al. Phys. Rev. Lett., 102, 254801 (2009).
Transverse slice emittance scaling with charge

LCLS results at 20 pC: slice emittance <0.2µm. Y. Ding et al., Phys. Rev. Lett. 102, 254801 (2009)

Emittance scaling. $\varepsilon_N$ in µm, Q nC
For $Q < 0.3$ nC the RF term is negligible. Ferrario et al., Nucl. Instr. And Meth. A57, 98 (2006).

$$\varepsilon_N = 1.4\sqrt{0.111Q^{2/3} + 0.18Q^{4/3} + 0.18Q^{8/3}}$$

The red dots are LCLS experimental results. The empirical factor 1.4 indicates a thermal emittance larger than theoretical value.
Another method to generate short pulses: the slotted emittance spoiler system.


Successfully used to generate few fs pulses at LCLS.

Courtesy P. Emma
Double X-Ray Pulses from a Double-Slotted Foil

Controlled time delay between x-ray pump and x-ray probe pulses.

2-Pulse Production with 2 slots

0.25 mm

0-6 mm

Pulses not coherent

Power (GW)

10

5

0

0-150 fs

2 fs

Time (fs)

Courtesy P. Emma

Challenges of 4th Generation Light Sources, C. Pellegrini
<table>
<thead>
<tr>
<th>What next?</th>
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<tbody>
<tr>
<td>Photon energy, keV</td>
</tr>
<tr>
<td>Pulse repetition rate, Hz</td>
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<tr>
<td>Pulse duration, fs</td>
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<tr>
<td>Coherence, transverse</td>
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<tr>
<td>Coherence, longitudinal</td>
</tr>
<tr>
<td>Coherent photons/pulse</td>
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<tr>
<td>Peak (Average) brightness, ph/s mm² mrad² 0.1% bandwidth</td>
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<td></td>
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<tr>
<td>Peak (Average) power, TW (kW)</td>
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<tr>
<td>Polarization</td>
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<tr>
<td>Multicolor pulses</td>
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</tbody>
</table>
Different, specialized FELs will be needed to satisfy these requests.

Example 1. An X-ray FEL oscillator is a very good candidate to produce a nearly transform limited pulse, with a line width a small as $10^{-6} - 10^{-7}$ (K.-J. Kim, Y. Shvydko, S. Reiche, Phys. Rev. Lett. 100, 244802 (2008)). Reaching the same line-width in an amplifier is practically impossible, even with seeding. The oscillator would generate a small number of coherent photons per pulse in a long pulses, 0.1 to 1 ps, at high, MHz, repetition rate, using a CW superconducting linac. The X-ray oscillator would use low emittance, low charge, $\sim 50$ pC, electron bunches. Challenges: low loss mirrors in the Ångstrom to nanometer region; high repetition rate, one to a few MHz, electron guns with the required emittance and linear longitudinal phase space distribution.
Wish list remarks

**Example 2.** Single molecule imaging requires a large number of photons, $>10^{13}$ in about 10 fs or less, $\sim 1$ TW peak power, leading to a very different FEL optimization. Measurements are done in a single shot, blowing up the sample. Matching the sample preparation time and the FEL repetition rate is important.

**Example 3.** Reaching into the femtosecond/attosecond region reduces the number of photons and the charge per bunch that is required. The sample is not destroyed by the X-ray pulse and the amount of data/shot is limited. A high FEL repetition rate, up to MHz, becomes very desirable. The beam energy can be reduced using the low emittance obtained at small charge/bunch.
Femtosecond X-ray protein nanochrystallography

\[ E_{\text{photon}} = 1.8 \text{ keV} \]
\[ N_{\text{photons}} \approx 10^{12}/\text{pulse} \]
\[ T_{\text{pulse}} \approx 70\text{fs} \]

a) Diffraction intensity with single 70 fs pulse; resolution 8.5 Å
b) pattern of the [001] zone obtained from merging data from >15,000 nano-crystal;
c) electron density of photosystem I from LCLS data and
d) from synchrotron data with a resolution of 8.5 Å.
Mimivirus

The 2nd experiment [M.M. Seibert et al., Nature 470, 78 (2011)] shows that high quality diffraction data can be obtained from a single X-ray pulse on a non-crystalline biological sample, a single mimivirus particle. 7Å, 70 fs, $10^{12}$ photons.
The nano-crystal imaging experiment used 70 fs long pulses of about $10^{12}$ photons of 1.8 keV. Resolution about 8.5 Å. Reducing the pulse duration to 10 fs or less, increasing the number of photons to $10^{13}$ and the energy to 8 keV, allows single shot measurements of single molecules with a few Å resolution, a great breakthrough.

To reach this goal we are studying the feasibility of a 1 TW, 10 fs X-ray FEL at 1.5 Å, using LCLS electron beam parameters.

Examples of TW FELs other applications: splitting the X-ray pulse for 3D imaging with multiple orthogonal beams; non-linear electrodynamics, like multi-photon creation of electron-positron pairs, imaging matter in disordered states.
TW feasibility studies

Existing hard X-ray FELs, like LCLS, operate in high gain SASE mode, starting from longitudinal density noise in the electron beam and reaching saturation. The saturation power is ~30-40GW.

Kroll, Rosenbluth and Morton [N.M. Kroll, P.L. Morton, and M.N. Rosenbluth, IEEE J. Quantum Electronics, QE-17, 1436 (1981)] proposed to increase the energy transfer from the electron to the photon beam by adjusting the undulator magnetic field to compensate for the electron energy losses, a “tapered” undulator.

A tapered undulator in combination with self-seeding can be used to reach the 1 TW level.
TW X-ray FELs: two studies for XFEL & LCLS


The system consists of a SASE amplifier, followed by a "self-seeding" crystal monochromator, and a long tapered undulator.

Results for LCLS and XFEL show that TW-level output power at 8 keV is feasible. For LCLS we have a total undulator length below 200 m for a 10 fs, 40 pC bunch charge, normalized transverse emittance 0.3-mm-mrad, peak current 4 kA, electron energy 13.6 GeV.
TW SELF-SEEDED FEL

Start with a SASE FEL, followed by a monochromator and a tapered undulator. Simulations for the LCLS-II variable gap undulator.

SASE undulator power ~20 GW. Spectrum ➔

Self-seeding with single crystal monochromator proposed by G. Geloni, V. Kocharyan and E. Saldin, DESY 10-053 (2010). It can also be done with a 4 crystals monochromator and/or a double (fresh) bunch scheme.
TW FEL: LCLS-II case

- 8.3 keV -- 1.5 Å (13.64 GeV)
- 40-pC; 4-kA peak current; 10 fs FWHM; emittance 0.3-µm
- Quadratic tapered field decreasing 13% to 200 m.
- \( \lambda_U = 3.2 \text{ cm}, \) undulator sections 3.4 m, 1 m break; \( < \beta > = 20 \text{ m} \)
- Longitudinal: close to transform limited
- X-ray pulse brightness increases by \( 10^3 \)

\[ \Delta \lambda/\lambda = 10^{-4} \text{ FWHM} \]
Large harmonics in the TW FEL

The bunching remains high in the long tapered undulator and harmonics are large. For a tapered planar undulator the third harmonic at 24 keV is about 100GW.

Fundamental (red); Second harm. (blue); Third harm. (green)
Some general considerations: fs and meV

X-ray FELs can be designed to generate very short, fs to as, pulse duration $\tau$, or to pulses with very small line-width and small photon energy $\Delta E_{ph}$. In the two cases the connection between charge and emittance leads to very different optimization.

$$Q_{Bunch} = \frac{IL_{Bunch}}{c}, \quad L_{Bunch} = \frac{\lambda}{(\Delta \lambda / \lambda)},$$

$$Q_{Bunch} = \frac{I \lambda^2}{c \Delta \lambda} = \frac{Ih}{\Delta E},$$

Short pulse $\rightarrow$ low charge

Small $\Delta E$ $\rightarrow$ high charge

Fs pulse duration and 1 kA current gives a charge of about 1pC and an emittance 0.03 $\mu$m.

At 1nm and $\Delta E$=1 meV, we need 3nC for a current of 1kA and the emittance is 2$\mu$m. CSR and other collective effects will strongly distort the beam phase space. A high gain amplifier becomes complicated and a low gain oscillator system should be considered.
Using the beam properties already obtained we can already reduce the size and cost of X-ray FELs. With new possible advances we will be able to produce electron beams with even higher phase-space density, and better characteristics as an FEL lasing medium.

- Lower emittance
- Lower beam energy for given $\lambda$
- Better compression
- Larger peak current
- Compact, cheaper FEL with same X-ray beam

Short list of new advances
Some new results and concepts to improve the “lasing medium”

- Seeding and self seeding
Some new results and concepts to improve the “lasing medium”

- Controlling the bunch current profile in a photo-injector to reduce wake-field effects M. Cornacchia et al., ST/F-TN-06/06
- Split compressor, J. Frisch and Y. Ding, private comm.
First step towards a TableTop FEL: undulator based diagnostic of LPA beams

- Uses combined gas jet + capillary discharge based LPA*
- Beam imaging with permanent quadrupoles onto undulator entrance ~4.5 meter away: observed first light at LOASIS/LBNL

XUV spectrometer  Magnetic spectrometer

~ 7.5 meter


Courtesy W. Leemans
Ultra-low emittance measured using x-ray spectroscopy of betatron radiation of laser plasma accelerator beams

- X-ray spectrum provides measure of beam size inside accelerator:
  - 0.1-0.15 micron
  - Simultaneous divergence and beam energy measurement provides normalized emittance

\[ \epsilon_x \approx \gamma \sigma_x \sigma_\theta \]

0.1-0.2 mm-mrad - normalized

Value consistent with simulations
Conclusions

1. X-ray FELs can be developed to fulfill most of our requirements: femto- to atto-second pulse duration, very small line width, ultra-high, TW, peak power, opening new windows in many areas of science.

2. Utilizing the extraordinary brightness of low-charge bunches it is possible to reduce the size and cost of the accelerator, particularly so for short pulses and using new short period undulators.

3. Longitudinal coherence can be pushed near the transform limit using single spike, self seeding, seeding, or an X-ray oscillator.

4. Developments in understanding beam physics and manipulation may enable more compact future hard x-ray FEL facilities.

5. Research on laser/plasma/wakefield accelerators and new electron sources can lead to compact, university scale, X-ray FELs.